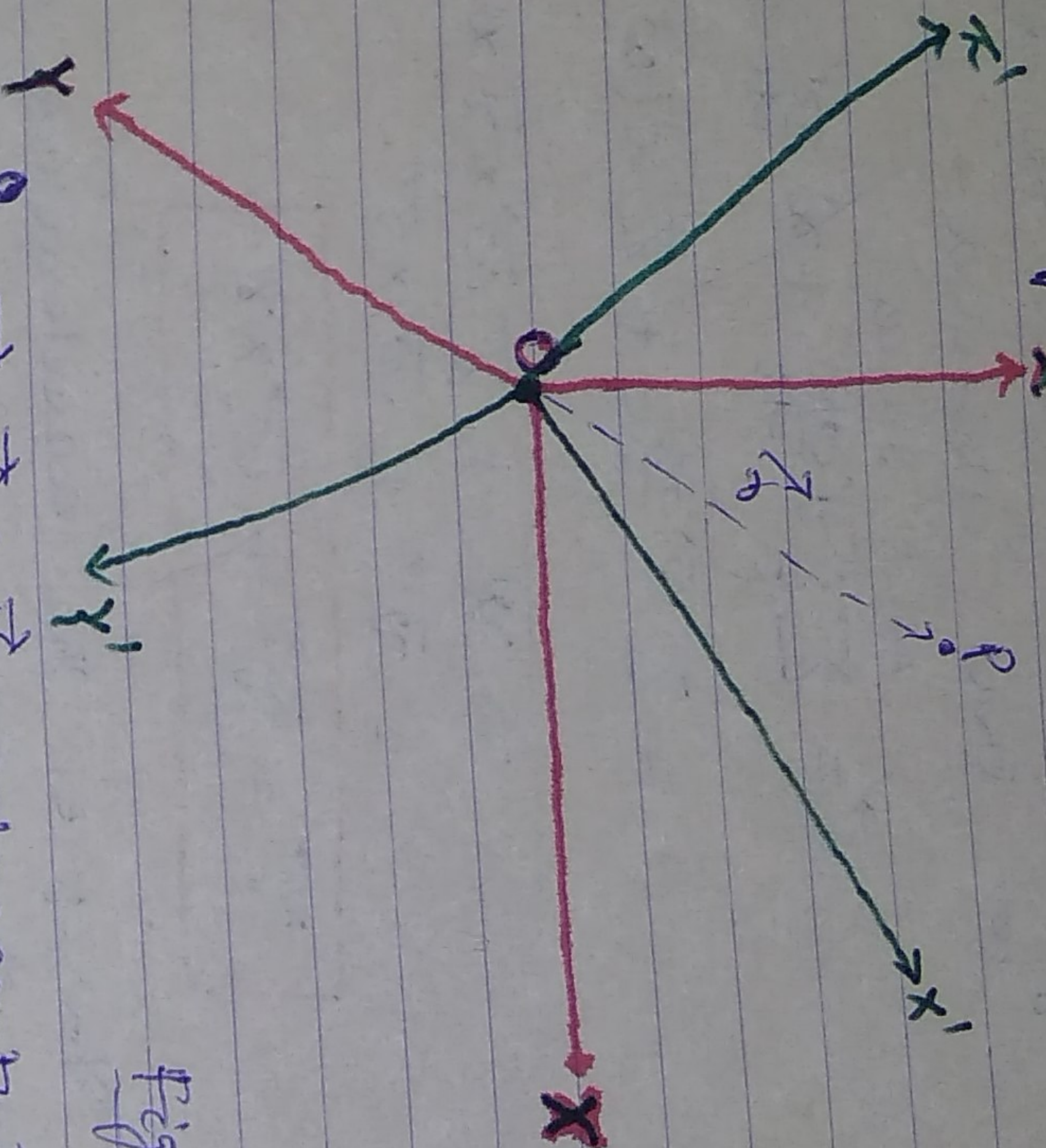


COROLLIS FORCE - ROTATING FRAME OF REFERENCE.

Let us consider a frame of reference x', y', z' rotating with a constant angular velocity " ω " with respect to a stationary frame of reference x, y, z , both frame having the same origin at O.



Figure

At t , $\vec{i}, \vec{j}, \vec{k}$ represents the unit vectors along x, y, z directions and $\vec{i}', \vec{j}', \vec{k}'$ the unit vectors along x', y', z' directions respectively. Then the position vector \vec{r} of a particle P will be given by,

(i) For x', y', z' system,
 $\vec{r} = \vec{i}'x' + \vec{j}'y' + \vec{k}'z'$ and

(ii) For x, y, z system,
 $\vec{r} = \vec{i}x + \vec{j}y + \vec{k}z$.

Therefore the velocity of the particle at P will be given by,

(i) For an observer in x', y', z' system,

$$\vec{v}' = \frac{d\vec{r}}{dt} = \sum \frac{d}{dt}(\vec{i}'x') = \sum \vec{i}' \frac{dx'}{dt}$$

And (ii) For an observer in x, y, z System,

$$\vec{v} = \frac{d\vec{r}}{dt} = \sum \frac{d}{dt} (\vec{r}' x')$$

Since the unit vectors \vec{i}', \vec{j}' & \vec{k}' are rotating for the observer in x, y, z System, they must be differentiable for that observer.

And Hence,

$$\vec{v} = \sum \left[\vec{r}' \frac{dx'}{dt} + x' \frac{d\vec{r}'}{dt} \right] \quad \left| \begin{array}{l} \because \vec{v} = \frac{d\vec{r}}{dt} \\ \vec{v} = \vec{\omega} \times \vec{r} \end{array} \right.$$

$$= \sum \vec{r}' \frac{dx'}{dt} + \sum x' (\vec{\omega} \times \vec{r}')$$

$$= \vec{v}' + \sum \vec{\omega} \times \vec{r}' x'$$

$$= \vec{v}' + \vec{\omega} \times \vec{r}$$

$$\boxed{\vec{v} = \vec{v}' + \vec{\omega} \times \vec{r}} \text{ --- (a)}$$

Therefore, the acceleration of the Particle for the observer in x, y, z System will be given by,

$$\frac{d\vec{v}}{dt} = \frac{d\vec{v}'}{dt} + \frac{d}{dt} (\vec{\omega} \times \vec{r})$$

$$= \frac{d}{dt} \sum \vec{r}' v'_x + \vec{\omega} \times \frac{d\vec{r}}{dt}$$

$$= \sum \vec{r}' \frac{d}{dt} v'_x + \sum v'_x \frac{d\vec{r}'}{dt} + \vec{\omega} \times \frac{d\vec{r}}{dt}$$

$$= \frac{d\vec{v}'}{dt} + \sum v'_x (\vec{\omega} \times \vec{r}') + \vec{\omega} \times \frac{d\vec{r}}{dt}$$

$$= \frac{d\vec{v}'}{dt} + \vec{\omega} \times \sum \vec{r}' v'_x + (\vec{\omega} \times \frac{d\vec{r}}{dt})$$

$$= \frac{d\vec{v}'}{dt} + (\vec{\omega} \times \vec{v}') + \vec{\omega} \times [\vec{v}' + (\vec{\omega} \times \vec{r})]$$

Using eqn (a)

$$\therefore \frac{d\vec{v}}{dt} = \frac{d\vec{v}'}{dt} + 2(\vec{\omega} \times \vec{v}') + \vec{\omega} \times (\vec{\omega} \times \vec{r}) \quad \text{--- (b)}$$

If m is the mass of the particle, then the force acting on it for the observer in x, y, z system.

$$\vec{F} = m \frac{d\vec{v}}{dt}$$

$$= m \frac{d\vec{v}'}{dt} + 2m(\vec{\omega} \times \vec{v}') + m\vec{\omega} \times (\vec{\omega} \times \vec{r})$$

$$= \vec{F}' + 2m(\vec{\omega} \times \vec{v}') + m\vec{\omega} \times (\vec{\omega} \times \vec{r})$$

$$\therefore \vec{F}' = \vec{F} - 2m(\vec{\omega} \times \vec{v}') - m\vec{\omega} \times (\vec{\omega} \times \vec{r}) \quad \text{--- (c)}$$

In this expression the term $[-2m(\vec{\omega} \times \vec{v}')] is called Coriolis force and the term $[-m\vec{\omega} \times (\vec{\omega} \times \vec{r})]$ is called Centrifugal force.$

These two forces are fictitious forces acting on the particle which arise due to the rotation of the frame of reference. If there is no rotation, i.e. if $\vec{\omega} = 0$, then the two forces will disappear.

CORIOLIS FORCE: \rightarrow The Coriolis force is therefore a fictitious force acting on a particle when the particle moves relative to a rotating frame of reference and which is mathematically defined by,

$$\vec{F}_c = -2m(\vec{\omega} \times \vec{v}')$$

where $\vec{\omega}$ is the angular velocity of the rotatory reference frame S' and \vec{v} is the velocity of the particle relative to the system S' .

The -ve sign in the expression \vec{F}_c indicates that the Coriolis force is in the direction opposite to that given by right hand screw rule. =

Qu 2. \rightarrow